

Inflationary Scenario in Bianchi Type II Space with Bulk Viscosity in General Relativity

Laxmi Poonia¹, Sanjay Sharma^{2*}

*Correspondence: sanjaysharma5416@gmail.com

¹⁻²Department of Mathematics and Statistics, Manipal University Jaipur, Jaipur, Rajasthan, India

ORCID:

Laxmi POONIA: <https://orcid.org/0000-0001->

Sanjay SHARMA: <https://orcid.org/0000-0001-9049-1349>

Abstract: This study is about bulk viscous inflationary model with flat potential under framework of LRS Bianchi type II metric. To drive relativistic solution of the field equations we choose the proportionally condition between coefficient of shear and expansion scalar which leads to a suitable relation $a = b^n$ between the metric functions where n is the constant other than one. Some dynamical features of the universe are also discussed.

Keywords: Bianchi Type II, Inflationary Model, Bulk Viscosity, Flat Potential, General Relativity

1. Introduction

In recent scenario, inflationary cosmology has become a curious subject of study to explaining the evolution of universe and formation of galactic structures, it is widely acceptable by many cosmologists that major cosmological problems like isotropy, homogeneity, flatness are successfully explained by the theory of inflation. The choice of anisotropic metric with system of field's equations allows us to construct a mathematical model of cosmos to understand the accelerated fate of the current physical universe. The system of field equations are basically set of non-linear differential equations we required solutions of it in various applications in astrophysics and cosmology. The analysis of microwave background has also provided some physical evidence regarding the cosmos inflation. Starobinsky [1] explained the model of initial universe after epoch. Guth [2] studies the various aspects of cosmos inflation by proposing the fact that false vacuum energy is responsible for this phenomenon. The role of Higgs fields with potential V are significantly used in various research. Many cosmologists [3-11] have derived different models to understand the theory of inflation and scalar field ϕ in various manners.

LRS Bianchi- II metric has a significant role in developing models which help to understand the early evolution of cosmos and inflationary nature in more details. Inflationary model within framework of bulk viscosity is very helpful to illustrate many physical and structural features in the dynamics of the current universe. Mishner [12] studied inflationary model under the effect of bulk viscosity in different manners. Heller and Klimek [13] has constructed a viscous fluid model without initial singularities. Gron [14-15] has discussed Bianchi - I space with bulk, nonlinear viscosity and in shearing mode. The significant role of bulk viscosity in cosmos inflation is studied by many researchers [16-18] and Bali et al. [19-21] in different aspects. Agrawal [22] investigate LRS

Bianchi II model inflationary model. Reddy [23] investigated Bianchi Type V space time for massless scalar field under flat potential. Sharma [24] derived Bianchi II inflationary model in general relativity. Bali and Poonia [25] constructed inflationary cosmological model under framework of Bianchi type III with bulk viscosity.

Motivated by this discussion, we have investigated Bulk viscous inflationary cosmological model with flat potential for LRS Bianchi type II metric. To get inflationary solution we suppose $\xi\theta = \alpha$ (constant) as proposed by Brevik et al.[26]. The paper work is classified as following sections given as: Section-2 concerned with metric and system of nonlinear fields equations. Section-3 contains solutions of fields equations in inflationary context. Section-4 are concerned with dynamical aspects of constructed model. Section-5 deals with conclusion and discussion.

2. Fields Equations with metric

LRS Bianchi-II space in orthonormal frame can be obtained by line element

$$ds^2 = g_{\mu\nu} \theta^\mu \theta^\nu, \quad g_{\mu\nu} = \text{diag} (1, 1, 1, -1) \quad (1)$$

Here $a(t)$ and $b(t)$ are the metric function

$$\text{where } \theta^1 = a(t)\omega^1, \theta^2 = b(t)\omega^2, \theta^3 = a(t)\omega^3, \theta^4 = dt$$

$$\text{with } \omega^1 = dx, \omega^2 = dy - xdz, \omega^3 = dz \quad (2)$$

from equation (1) and (2) metric can be written as

$$ds^2 = -dt^2 + a^2(t)dx^2 + b^2(t)(dy - xdz)^2 + a^2(t)dz^2 \quad (3)$$

The field of gravity minimally to scalar region with potential $V(\varphi)$ is given by

$$S = \int \left(R - \frac{1}{2} \varphi_{,i} \varphi_{,j} g^{ij} - V(\varphi) \right) \sqrt{-g} dx^4 \quad (4)$$

Einstein's field equation for the model is given by

$$R_i^j - \frac{1}{2} R g_i^j = -T_i^j \quad (5)$$

(In geometrical unit $8\pi G = c = 1$)

where energy momentum tensor for scalar field under consideration of bulk viscosity is given by

$$T_{ij} = \varphi_{,i} \varphi_{,j} - \left(\frac{1}{2} \varphi_{,l} \varphi^{,l} + V(\varphi) \right) g_{ij} - \xi \theta (g_{ij} + u_i u_j) \quad (6)$$

$$\text{with } \frac{1}{\sqrt{-g}} \partial_{,i} (\sqrt{-g} \varphi_{,i}) = -\frac{dV(\varphi)}{d\varphi} \quad (7)$$

where ξ and θ be the bulk viscosity coefficient and scalar expansion respectively

we assume comoving coordinate system as $u^i = (0,0,0,1)$

For LRS Bianchi type II metric (1), system of field equation (2) can be obtained as

$$2 \frac{a_{44}}{a} + \frac{a_4^2}{a^2} - \frac{3b^2}{4a^4} = -\left(\frac{1}{2} \varphi_4^2 - V(\varphi) - \xi \theta \right) \quad (8)$$

$$\frac{a_{44}}{a} + \frac{b_{44}}{b} + \frac{a_4 b_4}{ab} + \frac{1b^2}{4a^4} = -\left(\frac{1}{2} \varphi_4^2 - V(\varphi) - \xi \theta \right) \quad (9)$$

$$\frac{a_4^2}{a^2} + 2 \frac{a_4 b_4}{ab} - \frac{1}{4} \frac{b^2}{a^4} = \left(\frac{1}{2} \varphi_4^2 + V(\varphi) \right) \quad (10)$$

where indices 4 in metric coefficient shows the ordinary differentiation with respect to cosmic time

Equation (7) provides

$$\varphi_{44} + \left(2 \frac{a_4}{a} + \frac{b_4}{b} \right) \varphi_4 = - \frac{dV}{d\varphi} \quad (11)$$

The physical parameter of developed model which are significantly used to find solution of field equations and discussing geometrical features are given by

Proper volume of developed model (1) is given by

$$V = \sqrt{-g} = a^2 b \quad (12)$$

Expansion scalar (θ) for model is given by

$$\theta = u^i_{;i} = 2 \frac{a_4}{a} + \frac{b_4}{b} \quad (13)$$

Shear scalar for model is given by

$$\sigma^2 = \frac{1}{2} \sigma_{ij} \sigma^{ij} \quad (14)$$

$$\text{where } \sigma_{ij} = \frac{1}{2} (u_{i,k} D_j^k + u_{j,k} D_i^k) -$$

$$\frac{1}{3} \theta (g_{ij} + u_i u_j) \quad (15)$$

$$\text{and } D_i^j = \delta_i^j - u^i u_j$$

which leads to

$$\sigma^2 = \frac{1}{3} \left(\frac{a_4}{a} - \frac{b_4}{b} \right)^2 \quad (16)$$

Formula for Hubble parameter is given by

$$H = \frac{1}{3} \left(2 \frac{a_4}{a} + \frac{b_4}{b} \right) \quad (17)$$

we obtained deceleration parameter by given relation

$$q = - \frac{R_{44}/R}{R_4^2/R^2} \quad (18)$$

3. Fields Equations with solutions

Equation (8-10) is independent equations with five unknown a, b, φ, θ and ξ . For this purpose we required extra condition

$$\xi \theta = \alpha \text{ (Constant)} \quad (19)$$

i.e. coefficient of bulk viscosity is inversely proportional to expansion scalar and shear scalar (σ) is directly proportional to expansion (θ) which leads to

$$a = b^n n > 1 \quad (20)$$

we have also assumed constant potential V with flat region

$$\text{i.e.} \quad V(\varphi) = \zeta \text{ (Constant)}$$

Equation (11) leads to

$$\varphi_{44} + \left(2\frac{a_4}{a} + \frac{b_4}{b}\right)\varphi_4 = 0$$

on integrating we have

$$\varphi_4 = \frac{\varphi_0}{a^2 b} \quad (21)$$

From the equations (8) and (9), we obtained

$$\frac{a_{44}}{a} - \frac{b_{44}}{b} + \frac{a_4^2}{a^2} - \frac{a_4 b_4}{ab} - \frac{b^2}{a^4} = 0 \quad (22)$$

Equations (20) and (22) give

$$2b_{44} + 4n\frac{b_4^2}{b} = \frac{2}{n-1}b^{3-4n} \quad (23)$$

$$\text{we consider} \quad b_4 = f(b) \quad (24)$$

$$\text{which provide} \quad b_{44} = ff' \quad (25)$$

$$\text{where } f' = \frac{df}{db}$$

from equations (23),(24) and (25) we have

$$\frac{df^2}{db} + 4n\frac{f^2}{b} = \frac{2}{n-1}b^{3-4n} \quad (26)$$

which leads to

$$f^2 = \frac{1}{2(n-1)}b^{4-4n} + Db^{-4n} \quad (27)$$

where D is the integrating constant

From equation (17) and (18) we have

$$b_4 = \left[\frac{1}{2(n-1)}b^{4-4n} + Db^{-4n}\right]^{\frac{1}{2}} \quad (28)$$

Equation (28) leads to

$$\int \left[\frac{1}{2(n-1)}b^{4-4n} + Db^{-4n}\right]^{\frac{1}{2}} db = \pm(t - t_0)$$

where t_0 is the integration constant

The metric (1) can be reduces in to form

$$ds^2 = -\left[\frac{1}{2(n-1)}T^{4(1-n)} + DT^{-4n}\right]^{-1} dT^2 + T^{2n}[dX^2 + dZ^2] + T^2[dY - XdZ^2]^2 \quad (29)$$

using the transformation $b = T, x = X, y = Y, \text{ and } z = Z$

4. Structural and dynamical features of the model

The proper volume is given by

$$V = T^{2n+1} \quad (30)$$

The shear scalar (σ^2) of model is obtained by

$$\sigma^2 = \frac{1}{3}(n-1)^2 \left[\frac{1}{2(n-1)T^{2(2n-1)}} + D T^{-2(2n+1)} \right] \quad (31)$$

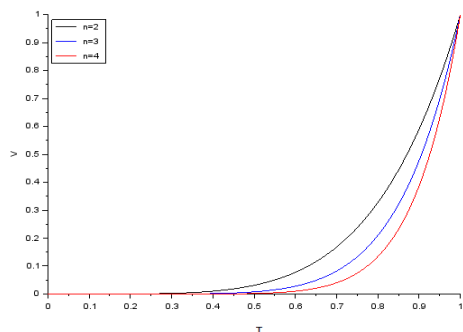


Fig. 1 volume (V) versus time (T)

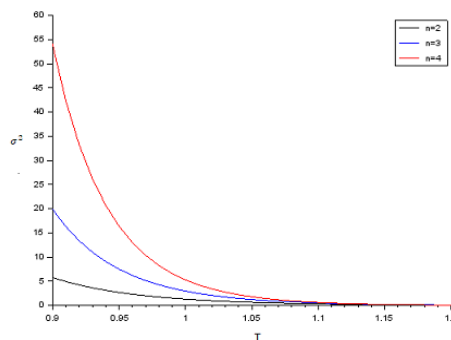


Fig. 2 coefficient of shear (σ) versus time (T)

The expansion (θ) for constructed model is given by

$$\theta = (2n+1) \left[\frac{1}{2(n-1)T^{2(2n-1)}} + D T^{-2(2n+1)} \right]^{\frac{1}{2}} \quad (32)$$

The result for hubble parameter (H) is given as

$$H = \frac{1}{3}(2n+1) \left[\frac{1}{2(n-1)T^{2(2n-1)}} + D T^{-2(2n+1)} \right]^{\frac{1}{2}} \quad (33)$$

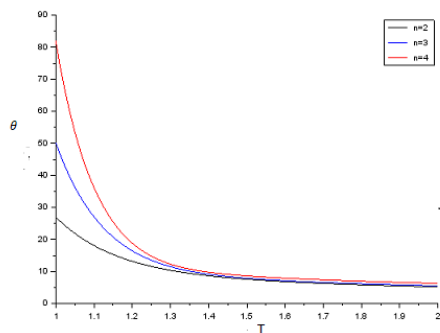


Fig. 3 coefficient of expansion (θ) versus time (T)

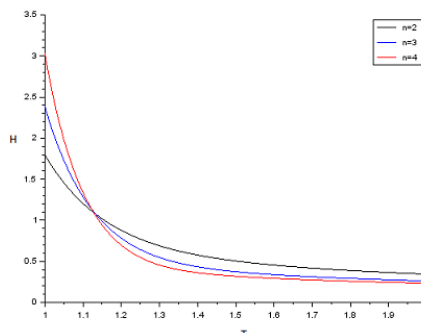


Fig. 4 Hubble parameter (H) versus time (T)

The deceleration parameter (q) for model is obtained by

$$q = -1 + \frac{3}{2n+1} \quad (33)$$

Coefficient of bulk viscosity (ξ) is given by

$$\xi = \frac{\alpha}{(2n+1) \left[\frac{1}{2(n-1)T^{2(2n-1)}} + D T^{-2(2n+1)} \right]^{\frac{1}{2}}} \quad (34)$$

Scalar field is given by

$$\varphi = \int \frac{\varphi_0}{T^{2n+1}} dT + C$$

5. Conclusion and discussion

The proper volume grows with time and become infinite at late time which indicates that cosmic inflation is possible in developed model. Since $\frac{\sigma}{\theta} = \text{constant}$ i.e. model maintained the anisotropy at late time, but the model become shear free and isotropic at $n=1$. The negative deceleration parameter shows accelerated phase of universe. The scalar of expansion and Hubble parameter become divergent at initial epoch, $t=0$ and tends to zero for infinite T and at $n = -\frac{1}{2}$ i.e. universe start with infinite expansion. The shear scalar is decreasing function of time and tends to zero for infinite large T . Higgs field decrease slowly and become finite for large time. The bulk viscosity coefficient leads to cosmic inflation in present scenario.

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